



1 **Intercomparison of Photoacoustic and Cavity Attenuated**
2 **Phase Shift Instruments: Laboratory Calibration and**
3 **Field Measurements**

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14 **Abstract:** The study of aerosol optical properties is essential to understand its impact on the global
15 climate. In our recent field measurement, a photoacoustic extincionometer (PAX) and a cavity attenuated
16 phase shift albedo monitor (CAPS-ALB) were used for online aerosol optical properties measurement.
17 Laboratory calibration with gas and particle samples were carried out to correct disagreements of field
18 measurements. During particle calibration, we adopted ammonium sulfate (AS) samples for scattering
19 calibration of nephelometer parts of both the instruments, then combined with number-size distribution
20 measurements into MIE model for calculating the value of the total scattering (extinction) coefficient.
21 During gas calibration, we employed high concentration NO₂ for absorption calibration of PAX resonator,
22 then further intercompared the extinction coefficient of CAPS-ALB with a cavity-enhanced spectrometer.
23 The correction coefficient obtained from the laboratory calibration experiments was employed on the
24 optical properties observed in the filed measurements correspondingly, and showed good result in
25 comparison with reconstructed extinction from the IMPROVE model. The intercomparison of the
26 calibrated optical properties of PAX and CAPS-ALB in field measurement were in good agreement with
27 slopes of 1.052, 1.024 and 1.046 for extinction, scattering and absorption respectively, which shows the
28 reliability of measurement results and verifies the correlation between the photoacoustic and the cavity
29 attenuated phase shift instruments.

30 **Key words:** Aerosol optical properties, instrument calibration, photoacoustic spectroscopy, cavity



31 attenuated phase shift spectroscopy

32 **1 Introduction**

33 Atmospheric aerosols can directly affect the earth's energy balance and cause global temperature
34 changes by absorbing and scattering solar radiation(Horvath, 1993;Haywood and Shine, 1995;Penner et
35 al., 2001). Therefore, considerable studies were undertaken to investigate the optical properties of aerosol
36 particles from different regions(Baynard et al., 2007;Petzold et al., 2013;Moosmüller et al., 1998). The
37 optical properties of regional aerosols depend on particle size distribution, mixing state and complex
38 refractive index, thus online measurements are necessary(Nakayama et al., 2015;Schwartz et al., 2010).
39 Furthermore, the calibration of instruments is a key step to ensure the reliability and quality of online
40 measurement data of aerosol optical properties.

41 Ideally, the complete set of aerosol optical properties are required measuring simultaneously,
42 including aerosol extinction, scattering and absorption coefficients, for aerosol optical closure studies.
43 The integrating nephelometry (IN) is an effective, economical and widely recognized method for online
44 obtaining aerosol scattering coefficient(Beuttell and Brewer, 1949;Heintzenberg and Charlson,
45 1996;Abu-Rahmah et al., 2006). Early on the systematic limitations of this technique were noted, that is
46 so-called truncation error caused by technically impossible to cover the full range of the scattering angle,
47 and which has mainly studied through numerical simulations with Mie model(S. Ensor and P. Waggoner,
48 1970;Anderson et al., 1996;Anderson and Ogren, 1998;Heintzenberg et al., 2006;Müller et al., 2009).
49 The measurement techniques for the extinction coefficient of atmospheric aerosols mainly include cavity
50 ring-down spectroscopy (CRDS) technique, cavity attenuation phase shift (CAPS) technique and cavity
51 enhanced absorption spectroscopy (CEAS) technique. CRDS has extremely high detection accuracy and
52 mature measurement system, which performed well in laboratory studies and field
53 measurements(O'Keefe and Deacon, 1988;Baynard et al., 2007;Berden et al., 2010;Pettersson et al.,
54 2004;Strawa et al., 2003). Related in its basic principle to CRDS, previously CAPS was used to calibrate
55 the reflectivity of mirrors also applied to measure atmospheric nitrogen dioxide(Kebabian et al., 2005;Ge
56 et al., 2013;Herbelin and McKay, 1981). It currently has been extended to the field of aerosol extinction
57 coefficient measurement(Kebabian et al., 2007;Petzold et al., 2013). Massoli et al. (2010) gave a detailed
58 description of CAPS results in the aerosol extinction coefficient measurements, including the first



59 laboratory characterization and field deployment. Onasch et al. (2015) calibrated the optical path length
60 error of CAPS with MIE model using monodisperse polystyrene spheres generated in the laboratory.
61 Rather than single wavelength measurements, CEAS with broadband light source applied for
62 atmospheric trace gas detection(Fiedler et al., 2003;Ball et al., 2004;Chen and Venables, 2011) was later
63 extended to quantitative aerosol extinction(Varma et al., 2013;Zhao et al., 2014;Suhail et al., 2019). The
64 filter-based methods are most commonly used for online measuring aerosol absorption
65 coefficient(Horvath, 1997;Hansen et al., 1982;Petzold and Schönlinner, 2004). Considering aerosol
66 morphology changes, multiple scattering and shielding effects, these methods require many correction
67 factors that limits the quality of measurement results(Bond et al., 1999;Weingartner et al., 2003).
68 Recently, the photoacoustic spectroscopy (PAS) technique(Terhune and Anderson, 1977;Bruce and
69 Pinnick, 1977;Adams et al., 1990), a direct method that can be easily-calibrated, has been developed into
70 a stable instrument in the field measurement of aerosol absorption(Moosmüller et al., 1998;Arnott et al.,
71 1999;Lack et al., 2006;Lewis et al., 2008;Sharma et al., 2013;Nakayama et al., 2015). Arnott et al. (2000)
72 calibrated their aerosol photoacoustic instrument by measuring the photoacoustic response in the
73 presence of NO₂ and compared its result with aethalometer. Lack et al. (2006) used ozone with a known
74 optical absorption level to calibrate the photoacoustic system with CRDS.

75 During our recent field campaign in Yangtze River Delta (YRD), the measurements of aerosol
76 optical properties showed discrepancies from different instruments, among which the extinction,
77 absorption, and scattering coefficients were measured by CAPS, PAS, and IN respectively(Du et al.,
78 2020). For investigation of the discrepancy between instruments and correction of the measurement data,
79 this study carried out an aerosol optical properties intercomparison measurement. During calibration
80 measurement, the extinction coefficient was calibrated with MIE model using mono-disperse particles
81 and the absorption coefficient was calibrated with transmission method using an absorbing gas, while
82 the scattering coefficient was calibrated with combination of above model and the method using no-
83 absorbing particles. In addition, an Incoherent Broad-Band Cavity Enhanced Absorption Spectroscopy
84 (IBBCEAS) setup was used to measure extinction coefficient of NO₂ for comparing with CAPS. Then
85 the correction factors obtained from the laboratory calibration experiments were employed on the data
86 observed in the filed measurement correspondingly and compared with the reconstructed extinction of
87 the interagency monitoring of protected visual environment (IMPROVE) model. Furthermore, the



88 calibrated field measurement results from photoacoustic and cavity attenuated phase shift instruments
89 were intercompared. For aerosol optical properties, different optical methods showed good agreement
90 and closure correlation after calibration, which has been rarely studied in laboratory calibration and field
91 measurement. In addition, the corrected field measurement data are more reliable for subsequent study
92 of aerosol optical properties in YRD region.

93 2 Materials and Methods

94 2.1 Instrument description

95 During calibration experiments, the optical properties of aerosol were measured by a Cavity
96 Attenuation Phase Shift-ALBedo monitor (CAPS-ALB) (Shoreline Science Reaserch, Japan) and a
97 Photoacoustic Extinctionmeter (PAX) (Droplet Measurement Technologies, US). In addition, a Scanning
98 Mobility Particle Sizer Spectrometer (SMPS) (Model 3938, TSI, US) was employed to measure the
99 number-size distribution for MIE model, and an IBBCEAS setup was used to measure NO₂ concentration
100 for extinction calculation. Above instrument details are summarized in Table 1.

101 **Table 1 Instrument Details**

Instrument	Parameters	Time resolution	Flow, Lmin ⁻¹	Wavelength, nm
CAPS-ALB	Extinction coefficient, Scattering coefficient [Mm ⁻¹]	1 s	0.85	530
PAX	Absorption coefficient, Scattering coefficient [Mm ⁻¹]	1 s	1	532
SMPS	Number size distribution [cm ⁻³]	5 min	0.3	-
IBBCEAS	NO ₂ concentration [ppb]	1 s	0.6	355-380

102 Aerosol sample flow was drawn into the PAX using an external vacuum pump, then split between
103 the wide-angle integrating reciprocal nephelometer and photoacoustic resonator for simultaneous online
104 measurements of light scattering coefficient and absorption coefficient. In the photoacoustic cavity, the
105 laser beam passing through the sample stream was modulated at the resonant frequency of the cavity, and
106 the light-absorbing molecules were heated and quickly transferred the heat to the receiving end of the
107 instrument, the pressure wave generated by periodic heating was detected by a sensitive microphone. The



108 calculation formula of absorption coefficient (b_{abs}^{obs}) is as follows (Rosenzweig, 1980):

$$109 \quad b_{abs}^{obs} = \frac{P_{mic} \cdot A_{res} \cdot \pi^2 \cdot f_{res}}{P_L \cdot (\gamma - 1) \cdot Q} \quad (1)$$

110 Where, P_{mic} is the pressure at the microphone at the resonant frequency f_{res} , P_L is laser power, A_{res} is the
111 geometric cross-section of the resonator, γ is the ratio of specific heat at constant pressure and volume,
112 Q is the quality factor of the resonator that calculated from temperature, pressure, and relative humidity
113 (RH).

114 The wide-angle integrating reciprocal nephelometer with a scattering integration angle of 6-174°
115 range used in PAX, which detects scattering light from a parallel beam through a cosine-weighted
116 detector. The detector located in the center of the cavity is fiber coupled to a photo-multiplier tube (PMT),
117 where the measured laser power is proportional to the total scattering cross section. The expression for
118 determining scattering coefficient (b_{sca}^{obs}) is given by (Abu-Rahmah et al., 2006):

$$119 \quad b_{sca}^{obs} = \frac{P_{PMT}}{P_L} \quad (2)$$

120 Where P_{PMT} is the value of the PMT signal with scattering background subtracted, P_L is measured laser
121 power. The scattering background was measured during the zeroing process of the instrument operation.

122 In addition, the extinction coefficient (b_{ext}^{obs}) considered as theoretical value in recommended
123 calibration method of PAX that can be obtained by measuring the intensity of transmitted light with a
124 photodetector combined with Lambert Beer's law as follow:

$$125 \quad b_{ext}^{obs} = \frac{\ln(I_0/I)}{L} \cdot 10^6 [Mm^{-1}] \quad (3)$$

126 Where, I_0 and I are the laser intensity with or without extinction substances, respectively. L is the path
127 length of the laser beam through the cavity in meters, here is 0.354 m. 10^6 is a conversion factor to
128 express extinction in Mm^{-1} .

129 The CAPS-ALB using an internal vacuum pump to introduce aerosol flow into the sample cell to
130 measure the extinction coefficient and scattering coefficient simultaneously. Nearly 1° truncation angle
131 integrating sphere integrating nephelometer (ISIN) has been employed in CAPS-ALB. The integrating
132 sphere with attached truncation reduction tubes located around the sample cell and PMT are equipped to
133 collect scattering light, which effectively reduces the angle truncation error (Varma et al., 2003). As a
134 typical kind of reciprocal nephelometer, its scattering coefficient (b_{sca}^{obs}) can also be calculated using Eq
135 (2).



136 The extinction measurement system of CAPS-ALB utilizing a visible light-emitting diode (LED)
137 with the luminescence as a light source and a sample cell incorporating two high reflectivity mirrors
138 centered at the wavelength of the LED and a vacuum photodiode detector. The extinction coefficient of
139 CAPS-ALB (b_{ext}^{obs}) is obtained by measuring the light attenuation of the visible long optical path with a
140 vacuum photodiode, and detecting the phase shift of the square wave frequency modulation heterodyne
141 detection of the light source, its expression as follow:

$$142 \quad b_{ext}^{obs} = (\cot \vartheta - \cot \vartheta_0) \cdot (2\pi f / c) \quad (4)$$

143 where \cot is the cotangent, c is the speed of light, f is the LED modulation frequency, T and P are the
144 sample temperature and pressure, respectively. The amount of phase shift (ϑ) is a function of fixed
145 instrument properties such as cell length, mirror reflectivity and modulation frequency, and of the
146 presence of aerosols (Kebabian et al., 2007). The term $\cot \vartheta_0$ is obtained from a periodic baseline
147 measurement (using particle-free air). It is worth mentioning that the effective optical-path error in the
148 sample cell of CAPS-ALB, which caused by the purge airflow of the mirror limits the space of the aerosol
149 samples, has been initially corrected in the internal calculation process. The original correction factor
150 was 0.7 that close to the value reported by Onasch et al. (2015), which generally calibrated with MIE
151 model calculation.

152 Our self-developed IBBCEAS device was used to measure gas concentration in the NO₂ comparison
153 experiment (Chen and Venables, 2011). The IBBCEAS measures the light intensity change of the light
154 source through the optical cavity, then inverts the concentration of the gaseous samples. When a pair of
155 high-reflectivity plano-concave mirrors with a reflectivity of R are composed of an optical cavity with a
156 length of L that is illuminated by continuous broadband incoherent light, the output light intensity I is
157 equal to the sum of the output light intensity of each order. Combined with Lambert Beer's law, the
158 expression for extinction coefficient $b_{ext-CEAS}$ at measured wavelength as follow (Fiedler et al., 2003; Ball
159 et al., 2004):

$$160 \quad b_{ext-CEAS}(\lambda) = \left(\frac{I_0(\lambda)}{I(\lambda)} - 1 \right) \left(\frac{1 - R(\lambda)}{L} \right) = \Sigma \sigma_i(\lambda) N_i \quad (5)$$

161 Here, I_0 is light intensity without absorbing matter, σ_i and N_i are absolute extinction cross section
162 and concentration of species i . I_0 , I , R , σ_i and N_i are functions of wavelength. Therefore, for different
163 detection wavelengths, the extinction coefficient cannot be compared directly. A simple method is to
164 establish a relationship with the species concentration. Fitting the extinction cross-section σ_i to the



165 extinction coefficient ($b_{ext-CEAS}$), the concentration of the measured gas N_i can be inverted. Noting that the
166 reflectivity R of the cavity mirrors in IBBCEAS has been calibrated before our experiments, so the result
167 of IBBCEAS can be considered as absolute value.

168 The number size distribution for MIE model calculation was obtained from SMPS, which consists
169 of an Electrostatic Classifier (Model3082, TSI, US) and a Condensation Particle Counter (CPC) (Model
170 3750, TSI, US). The Electrostatic Classifier was used with a Long Differential Mobility Analyzer (LONG
171 DMA) (Model 3081, TSI, US), its particle-size selection range is 14.1~736.5 nm, with a sample flow of
172 0.3 L min^{-1} and a sheath flow of 3 L min^{-1} . The aerosol sample passes through the radioactive neutralizer
173 to be charged, then enters the DMA to select particles of different particle sizes by changing the voltage.
174 The number of selected particles is counted to after the process of hygroscopic growth in CPC, which
175 has an uncertainty of within $\pm 10 \%$ in measuring particle concentration (Petzold et al., 2013).

176 2.2 Experimental

177 Based on the above principles, we adopted the following experimental procedures to compare PAX
178 and CAPS-ALB as Fig. 1 shows. The blue solid line represents the process of particle calibration, the red
179 solid line represents the procedures of gas calibration. All joints have been leak tested to ensure tightness.

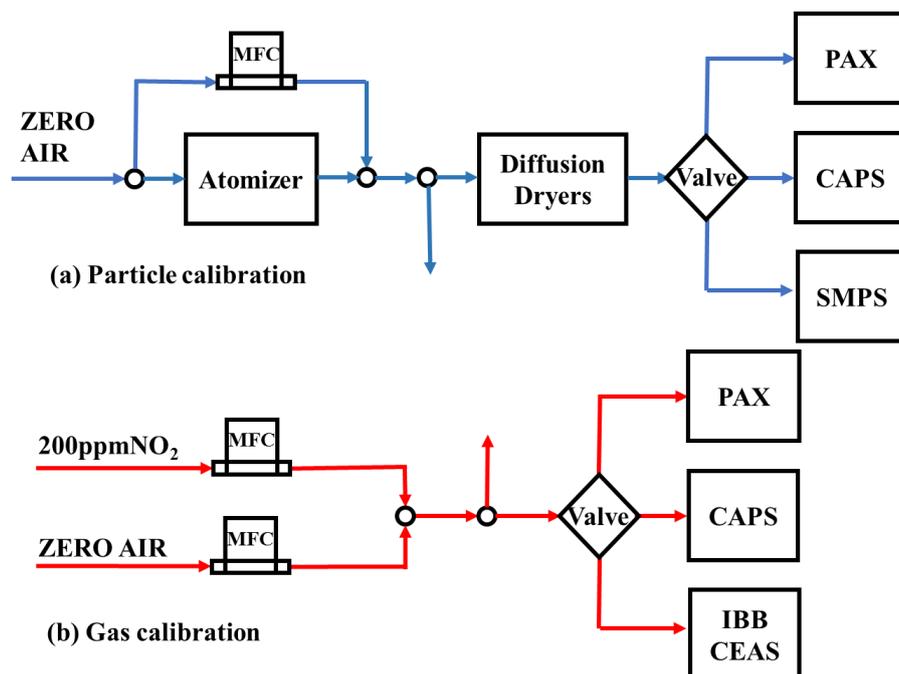


Figure 1: Experimental schematics (a) Particle calibration (b) Gas calibration.

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181



182 a. Particle calibration

183 For systematic errors, such as angle truncation, laboratory generated nebulized ammonium sulfate
184 (AS) (AR 99%, Aladdin Chemical) aerosols were used to calibrate and test the nephelometers of CAPS-
185 ALB and PAX using same experimental procedure as follow. AS aqueous solution was nebulized by an
186 atomizer (Model 9032, TSI, US) with filter air at a constant inlet pressure of 20 psi, which can generate
187 a stable outlet flow rate of $\sim 5 \text{ L min}^{-1}$. As shown by the blue solid line in the Fig. 1(a), the nebulized
188 aerosol flow was diluted with filtered air to adjust its concentration, and then dried using diffusion dryers
189 with silica gel that reduced the sample *RH* to $\sim 10 \%$ before delivery to the instruments, where the excess
190 airflow was discharged by bypass. Only opening PAX or CAPS-ALB valves, the dry aerosol flow was
191 connected to the instruments sampling port for at least $\sim 5 \text{ min}$ until the measured value stabilizes. The
192 entire flow system used conductive silicone tubing's and reduced bending to minimize the loss of
193 particles during aerosol transportation. For high-concentration of non-absorbing AS aerosol with
194 refractive index of $1.53+0.00i$, the absorption effect can be ignored. Therefore, the scattering calibration
195 factor (f_{sca}^{obs}) was calculated by comparing the measured extinction coefficient (b_{ext}^{obs}) and scattering
196 coefficient (b_{sca}^{obs}) (Lewis et al., 2008; Cross et al., 2010).

197 For the purpose of estimating the scattering or extinction coefficients measured in above
198 experiments and further correcting the absolute total scattering (extinction) coefficient, we performed an
199 additional calibration using polystyrene latex (PSL) spheres with Mie model. This model is a rigorous
200 analytical solution of the scattering field distribution of monochromatic light illuminates on spherical
201 particles (Born and Wolf, 1999). Thus, assuming the particles to be round, it is considered feasible to
202 apply Mie model to retrieve the number size distribution for calculating the total scattering coefficient of
203 atmosphere aerosol. The scattering and extinction efficiency factor Q_{sca} and Q_{ext} can be calculated from
204 the function of particle complex refractive index, light source wavelength and size distribution (Wu et
205 al., 2018; Bohren and Huffman, 1983). By integrating the particle cross-sectional area $\pi D^2/4$, particle
206 number concentration $N(D)$, and $Q_{sca/ext}$ on the particle diameter D distribution, yields the calculated
207 scattering and extinction coefficient $b_{sca/ext}^{MIE}$ as follow expression:

$$b_{sca}^{MIE} = \int_0^{\infty} Q_{sca} \cdot \frac{\pi D^2}{4} \cdot N(D) \cdot dD \quad (6)$$

$$b_{ext}^{MIE} = \int_0^{\infty} Q_{ext} \cdot \frac{\pi D^2}{4} \cdot N(D) \cdot dD \quad (7)$$



209 The experiments incorporating mono-disperse PSL spheres with complex refractive index
210 $1.60+0.00i$ and diameter of 350 ± 6 nm (Thermo Scientific) were carried out follow the calibration
211 procedures of scattering calibration (replace PSL with AS). Opened the SMPS valve and connected the
212 diluted dry aerosol flow to the its sampling port, then continuously measured together with CAPS-ALB
213 or PAX for ~ 20 min for collecting at least three sets of effective data of particle-size distribution at each
214 concentration. By comparing the MIE calculated average with the measured value for multiple
215 concentrations, the MIE model correction factor ($f_{sca}^{PAX-MIE}, f_{ext}^{CAPS-MIE}$) can be determined.

216 b. Gas calibration

217 As noted in previous studies (Arnott et al., 2000), the PAS resonator acoustic calibration used
218 sufficiently high concentrations of absorbing gas to generate a huge absorption, so that the Rayleigh
219 scattering was negligible. Therefore, our experiment adopted high concentration NO_2 for absorption
220 calibration and determined the absorption correction factor (f_{abs}^{obs}) from comparison of measured
221 absorption (b_{abs}^{obs}) and extinction (b_{ext}^{obs}) coefficients without knowing NO_2 concentration. In the case
222 of only PAX valve opened, by diluting 200 ppm NO_2 in different dilution ratios, the filtered air and NO_2
223 mixture were introduced to PAX for ~ 5 min, in which the flow of filtered air and NO_2 were controlled
224 by the mass flow controller to specified proportion, respectively. The entire flow system used Teflon
225 tubes to minimize NO_2 loss and contaminations, and a bypass was set to ensure the stability of the sample
226 flow and pressure.

227 Subsequently, considering the possible particulate loss of CAPS-ALB calibration, IBBCEAS and
228 CAPS-ALB was used to measure NO_2 samples simultaneously for comparing the measured extinction
229 coefficient in gaseous way. This experiment was carried out based on the experimental procedure for
230 PAX absorption calibration, though closing PAX route and simultaneously opening the valves of CAPS
231 and IBBCEAS. Based on the limitation of IBBCEAS the NO_2 concentration was controlled below 1 ppm
232 and each concentration was maintained for at least about 15 min until the measured value stabilizes. By
233 multiplying the NO_2 concentration measured by IBBCEAS and the NO_2 extinction cross section from
234 previous study of Voigt et al. (2002) at the CAPS-ALB detection wavelength (530nm), the conversion
235 result of the extinction coefficient ($b_{ext-CEAS}$) measured by IBBCEAS was obtained. Thus, the extinction
236 correction factor ($f_{ext}^{CAPS-CEAS}$) from comparison with IBBCEAS can be determined.

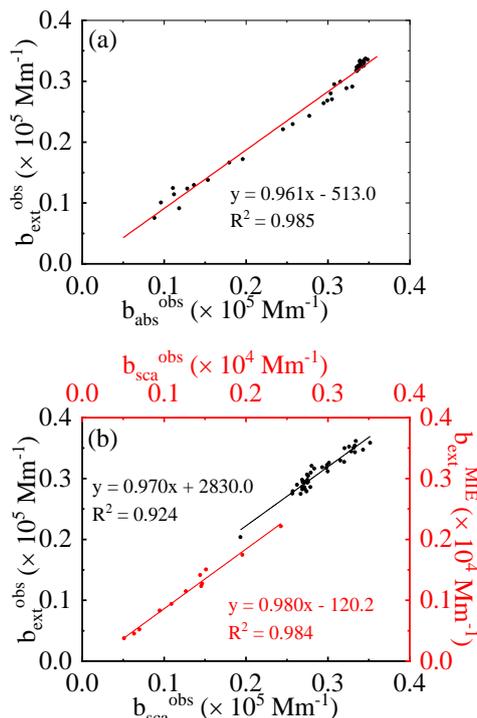


237 **3. Results and Discussion**

238 **3.1 PAX calibration results**

239 In our calibration experiments for PAX, with assumption of linearity in calibration down to the
 240 detection limit of the instruments (Arnott et al., 2000), the high concentration of absorbing gas and
 241 scattering particles generated a huge absorption and scattering effect that weakened the interference of noise
 242 for corrected the response curve of the PAX photoacoustic resonator and nephelometer respectively.

243 Fig. 2(a) shows the relationship between the absorption coefficient ($b_{\text{abs}}^{\text{obs}}$) and the extinction
 244 coefficient ($b_{\text{ext}}^{\text{obs}}$) in the NO_2 measurement results. The slope of fitted line in Fig. 2(a), which represent
 245 $f_{\text{abs}}^{\text{obs}}$, were determined to be 0.961 ± 0.019 with correlation factor $R^2 \sim 0.985$. The calibration result
 246 showed that the absorption measurement of PAX only needs slight correction and has high accuracy. Fig.
 247 2(b) presents typical correlation plots comparing the extinction coefficient from transmissivity ($b_{\text{ext}}^{\text{obs}}$)
 248 for AS samples (The black solid dot) and MIE model calculation ($b_{\text{ext}}^{\text{MIE}}$) for PSL samples (The red solid
 249 dot) with the measured scattering coefficient ($b_{\text{sca}}^{\text{obs}}$) respectively in PAX scattering calibration, where
 250 the extinction and scattering are theoretically equivalent due to negligible absorption.



251
 252

Figure 2: PAX calibration results: (a) Comparison of the measured extinction and absorption



253 **coefficient. (b) Comparison of the measured and MIE-model calculated extinction coefficient with the**
254 **measured scattering coefficient**

255 In Fig. 2(b), the slope of the black solid line indicates the measured scattering correction factor
256 (f_{sca}^{obs}) that was determined to be 0.970 ± 0.046 with correlation factor $R^2 \sim 0.924$. Moreover, we calculated
257 the absolute extinction coefficient with MIE model for further correction. Here, limited by the detection
258 range, another set of coordinate system was used for comparison. The slope of the red solid line that
259 indicates the MIE model scattering correction factor ($f_{sca}^{PAX-MIE}$) were determined to be 0.980 ± 0.039 with
260 correlation factor $R^2 \sim 0.984$. The scattering correction factors from transmission method and MIE model
261 were within acceptable range of the truncation error, and had only $\sim 1\%$ discrepancy in different
262 measurement range, showing that a good agreement between the two methods and the reliability of PAX
263 scattering calibration result.

264 **3.2 CAPS-ALB calibration results**

265 In the CAPS-ALB calibration experiment, we first utilized PSL spheres to correct its extinction
266 coefficient through MIE model calculation, and then employed AS samples to correct its scattering
267 coefficient comparing the calibrated extinction coefficient. In addition, we used self-developed
268 experimental IBBCEAS device for further verifying the correction factor calculated by MIE model.

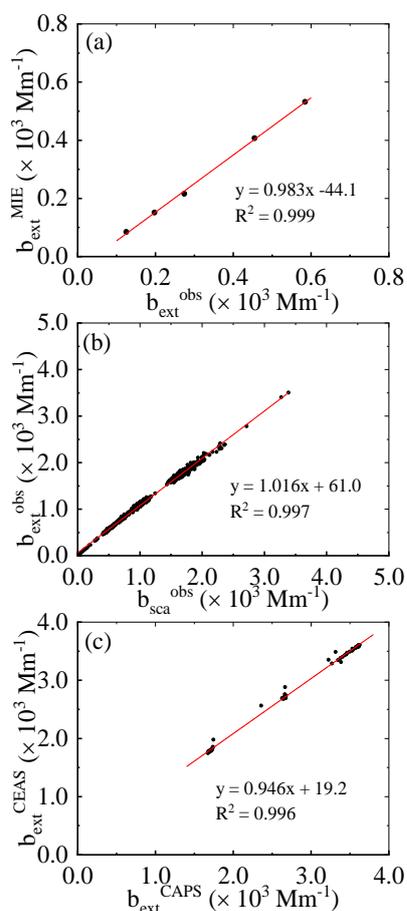
269 Fig. 3(a) shows correlation of extinction measured by CAPS-ALB (b_{ext}^{obs}) and extinction calculated
270 by MIE model (b_{ext}^{MIE}) for 350nm mono-disperse PSL spheres. The slope in Fig. 3(a) represents
271 extinction correction factor ($f_{ext}^{CAPS-MIE}$) were determined to be 0.983 ± 0.018 with correlation factor
272 $R^2 \sim 0.999$. It shows that the good accuracy of original calibration factor for the effective optical path error,
273 only slight adjustment was required. The other factors that might affect the extinction calibration is the
274 uncertainty of the aging effects of LED and detectors (PMT and vacuum photodiode), which has different
275 effects according to cell geometry.

276 Correlation plots comparing scattering coefficient (b_{sca}^{obs}) and extinction coefficient (b_{ext}^{CAPS}) for
277 AS samples measured by CAPS-ALB are shown in Fig. 3(b). According to its linear fitting result, f_{sca}^{obs}
278 were determined to be 1.016 ± 0.002 with correlation factor $R^2 \sim 0.996$. It showed that the measured
279 scattering coefficient has high accuracy, and verified the reliability of extinction correction factor of
280 CAPS-ALB ($f_{sca}^{CAPS-MIE}$).

281 Fig. 3(c) presents the comparison between measured extinction coefficient of CAPS (b_{ext}^{CAPS}) and



282 IBBCEAS (b_{ext}^{CEAS}) for NO₂ samples. The slope of Fig.3(c) represents extinction correction factor
283 ($f_{ext}^{CAPS-CEAS}$) were determined to be 0.946 ± 0.007 with correlation factor $R^2 \sim 0.998$. The experimental
284 correction factor of IBBCEAS ($f_{ext}^{CAPS-CEAS}$) was consistent with the theoretical correction factor of the
285 MIE model ($f_{ext}^{CAPS-MIE}$) within an acceptable error range of 4 %, proving that the reliability of MIE model
286 calculation and the applicability of CAPS-ALB calibration, no matter whether choosing gas or particle
287 ways.



288
289 **Figure 3: CAPS-ALB calibration results: (a) Comparison of the measured and MIE-calculated**
290 **extinction coefficient. (b) Comparison of the measured extinction and scattering coefficient. (c)**
291 **Comparison between measured extinction coefficient of CAPS-ALB and IBBCEAS**

292 3.3 Calibrated field measurement

293 The field measurements were carried out in the Gehu area of southwest Changzhou City, Jiangsu



294 Province (31°63' N, 119°90' E) from 25 May to 27 June before the rainy season in 2019. Changzhou has
295 a location in the center of the Yangtze River Delta and has a subtropical monsoon climate. The
296 measurement site was surrounded by 60 % of ecological wetlands and green gardens, and 20 % of
297 territorial waters, which results represented the regional ambient conditions of the Yangtze River Delta
298 before the rainy season. The sampling point was located on the top floor of a building at the height of 15
299 m above ground and all sampling tubes used a cyclone size cutter (URG, 2.5 μm, 5 lpm).

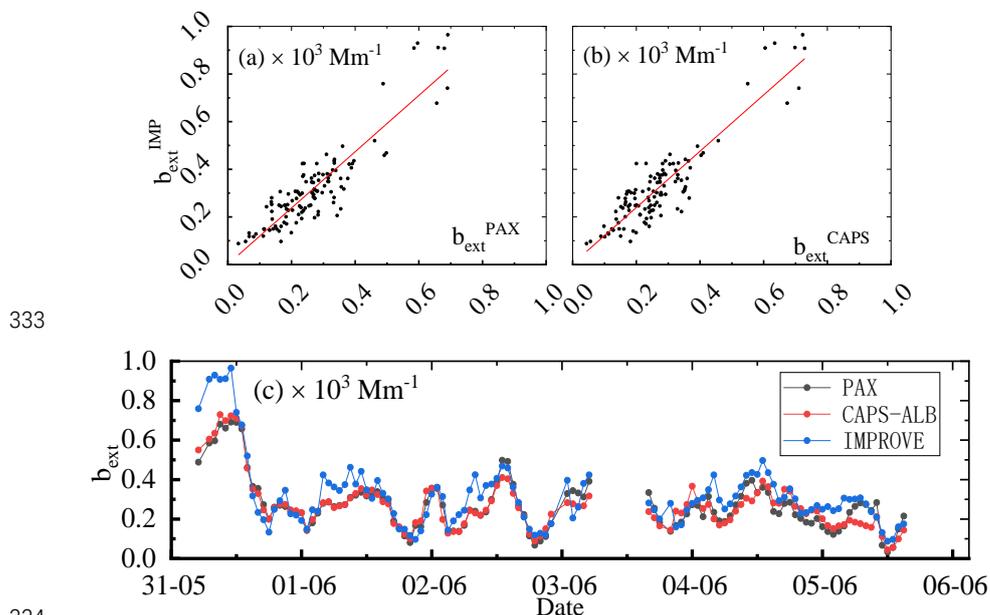
300 The correction factor obtained from the laboratory calibration experiments was employed on the
301 optical properties observed in the filed measurement correspondingly. For comparison, the IMPROVE
302 model was applied to identify major chemical components contributing to light extinction during field
303 measurement (Pitchford et al., 2007; Tao et al., 2014), where major chemical components including
304 water-soluble inorganic ions, organic carbon (OC), and elemental carbon (EC) were analyzed and
305 quantified by a monitor for aerosols and gases in ambient air (MARGA) (ADI 2080, Metrohm,
306 Switzerland) and an OC/EC analyzer (Model RT-4, Sunset, US). The simplified general formula of
307 IMPROVE model used in reconstruction of total scattering (extinction) coefficient (b_{ext}^{IMP}) can be
308 expressed as (Xia et al., 2017):

$$\begin{aligned} 309 \quad b_{ext}^{IMP} = & 2.2 \times f_s(RH) \times [Small (NH_4)_2SO_4] + 4.8 \times f_L(RH) \times [Large (NH_4)_2SO_4] \\ 310 \quad & + 2.4 \times f(RH) \times [Small NH_4NO_3] + 5.1 \times f(RH) \times [Large NH_4NO_3] \\ 311 \quad & + 2.8 \times [Small OM] + 6.1 \times [Large OM] + 1.7 \times f_{ss} \times [SS] + 1.0 \times [FS] \\ 312 \quad & + 0.6 \times [CM] + 8.28 \times [EC] \\ 313 \quad & [Large X] = [Total X]^2/20, [Total X] < 20, \\ 314 \quad & [Large X] = [Total X], [Total X] \geq 20, \\ 315 \quad & [Small X] = [Total X] - [Large X] \end{aligned} \quad (8)$$

316 where [X] represent the mass concentration of aerosol chemical component X, μg/m³; Ammonium
317 Sulfate [(NH₄)₂SO₄] = 1.375 [SO₄²⁻]; Ammonium Nitrate [NH₄NO₃] = 1.29 [NO₃⁻]; Organic Matters
318 [OM] = 1.6 [OC]; Sea Salt [SS] = 1.8 [Cl⁻]; Fine Soil [FS] = 2.2 [Al] + 2.49 [Si] + 1.94 [Ti] + 1.63 [Ca]
319 + 2.42 [Fe]; Coarse Mass [CM] = [PM10] - [PM2.5]; $f_s(RH)$, $f_L(RH)$ and f_{ss} represent RH growth curves
320 of sulfate, nitrate, and SS (Jung et al., 2009). Due to the lack of soil element information, Ca²⁺ was
321 assumed to account for 5 % of the concentration of fine soil mass based on previous studies, thus [FS] =
322 20[Ca²⁺] (Amato and Hopke, 2012).



323 Considering unavailable period of aerosol composition measurement (Due to status of MARGA),
324 only from 1 to 6 June were selected for the comparison. Fig. 4 (a) and (b) showed intercomparison of the
325 measured extinction coefficient of PAX and CAPS-ALB with IMPROVE-calculated extinction
326 coefficient, the linear fitting slopes are 1.182 and 1.183 with the correlation factor R^2 of 0.807 and 0.824,
327 respectively. Comparing the correlation factor, it is in good agreement with Shanghai (0.83) and
328 Hangzhou (0.81) in previous studies (Wang et al., 2016). Thus, it can be concluded that the IMPROVE
329 model has good applicability in Gehu area. Here, the extinction of PAX was the sum of the measured
330 absorption and scattering. In addition, Fig. 4(c) showed a timing diagram of the extinction coefficient
331 from PAX, CAPS-ALB measurement and IMPROVE model calculation. It showed a good agreement
332 between the measured and theoretical value and proved the reliability of our measurement data.



334
335 **Figure 4: Intercomparison of the measured extinction coefficient of (a) PAX and (b) CAPS-ALB with**
336 **IMPROVE-calculated extinction coefficient during field measurement (1-6 June 2019), and (c) the timing**
337 **diagram of the extinction coefficient from PAX, CAPS-ALB measurement and IMPROVE model**
338 **calculation.**

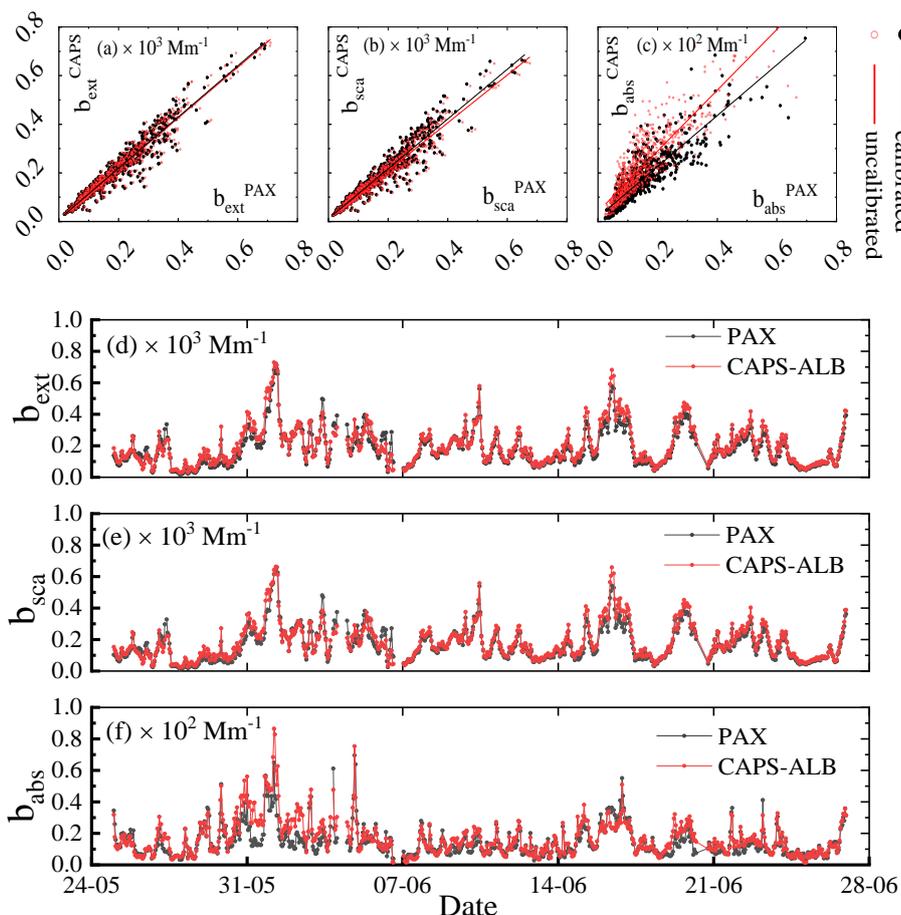
339 Then the CAPS-ALB and PAX corresponding optical properties of field measurement were
340 compared respectively in the case of calibrated and uncalibrated as Fig. 5(a), (b) and (c) showed. Here,
341 the extinction coefficient of PAX has been mentioned above as well as the absorption of CAPS-ALB was
342 the difference between the measured extinction and scattering. The linear fitting slope was 1.052, 1.024



343 and 1.046 from comparison of PAX and CAPS calibrated extinction, scattering and absorption coefficient,
344 with the correlation factor R^2 as 0.936, 0.924 and 0.772. Comparing the calibrated and uncalibrated
345 results, only slight corrections existed in the extinction and scattering coefficients, while the discrepancy
346 in the absorption coefficient has been corrected from ~30 % to less than 5 %. It can be considered that
347 the optical properties measured from PAX and CAPS-ALB with different measurement principles had a
348 good agreement, which in turn proved the reliability of our laboratory calibration results and the closure
349 correlation of CAPS-ALB and PAX measurements.

350 In addition, through deleting the time points of instruments data under zero calibration and abnormal
351 working conditions, the overall trend of calibrated extinction, scattering and absorption coefficients
352 during the measurement period (from 25 May to 27 June) were obtained as shown in Fig. 5(d), (e) and
353 (f). For the aerosol optical properties of the measurement region, it showed a dominated contribution of
354 scattering effect to the extinction coefficient, and a low levels of absorption coefficient.

355 The different internal structure of the nephelometers, even using the same principle, caused the
356 slight difference in the measured scattering coefficient. While the relatively small amount of absorption
357 coefficient of CAPS-ALB derived from extinction subtracted scattering coefficient has been greatly
358 affected. Therefore, the absorption coefficient which is difficult to quantify, was verified by CAPS-ALB
359 via correcting the scattering coefficient and the relationship of optical properties.



360 **Figure 5: Intercomparison of the CAPS-ALB and PAX (in the case of calibrated and uncalibrated) for (a)**
 361 **extinction, (b) scattering and (c) absorption coefficients during field measurement (From 25 May to 27**
 362 **June), and the timing diagram of the calibrated (d)extinction, (e)scattering and (f)absorption coefficients of**
 363 **CAPS-ALB and PAX.**

364 4. Conclusion

365 In this work we carried out aerosol optical properties inter-comparison measurements from
 366 photoacoustic and cavity attenuated phase shift instruments. The instruments were calibrated via
 367 laboratory experiments and the corrected field measurement data have also been intercompared. Thus,
 368 following points can be concluded:

- 369 (1) The laboratory results showed that disagreements exist between the two instruments before
 370 calibration. The scattering coefficient part plays a crucial role as the bridge in constructing correlation of



371 both instruments. Then the corrected extinction and absorption coefficients from both instruments were
372 intercompared well.

373 (2) The intercomparison of calibrated absorption and extinction coefficients in a field measurement
374 using photoacoustic and cavity attenuated phase shift instruments showed good agreement. Therefore,
375 laboratory calibrations were used for corrections for ensuring the quality of field data and further analysis
376 of radiative study.

377 **Data availability**

378 The raw data from the experiments are available upon request (j.chen@usst.edu.cn).

379 **Author Contribution**

380 **Jialuo Zhang:** Data curation, Methodology, Formal analysis, Visualization, Writing- Original draft
381 preparation. **Jun Chen:** Conceptualization, Investigation, Methodology, Supervision, Funding
382 acquisition, Writing-Review and Editing. **Meng Wang:** Validation. **Mingxu Su:** Supervision. **Wu Zhou:**
383 supervision. **Ravi Varma:** Investigation, Writing- Review and Editing, Methodology. **Shengrong Lou:**
384 Writing - Review and Editing, Validation, Funding acquisition.

385 **Competing Interests**

386 The authors declare that they have no conflict of interest.

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